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Dopant-Free Hole Transport Materials Afford Efficient and Stable Inorganic Perovskite Solar Cells and Modules

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Abstract: The emerging CsPbI₃ perovskites are highly efficient and thermally stable materials for wide-band gap perovskite solar cells (PSCs), but the doped hole transport materials (HTMs) accelerate the undesirable phase transition of CsPbI₃ in ambient. Herein, a dopant-free D- π -A type HTM named CI-TTIN-2F has been developed which overcomes this problem. The suitable optoelectronic properties and energy-level alignment endow CI-TTIN-2F with excellent charge collection properties. Moreover, CI-TTIN-2F provides multisite defecthealing effects on the defective sites of CsPbI₃ surface. Inorganic CsPbI₃ PSCs with CI-TTIN-2F HTM feature high efficiencies up to 15.9%, along with 86% efficiency retention after 1000 h under ambient conditions. Inorganic perovskite solar modules were also fabricated that exhibiting an efficiency of 11.0 % with a record area of 27 cm². This work confirms that using efficient dopant-free HTMs is an attractive strategy to stabilize inorganic PSCs for their future scale-up.

Introduction

Hybrid organic-inorganic perovskite solar cells (PSCs) have demonstrated remarkable progress in power conversion efficiencies (PCEs) from 3.8% to 25.5% in the past several years, showing their great potential in next-generation low-cost photovoltaic technology.^[1] However, the poor thermal stability of hybrid organic-inorganic perovskite materials hinders their large-scale commercialization.^[2] Thus, a signifi-

cant part of the research focus has been shifted to the allinorganic perovskite structures without volatile organic components.^[3] Among the various inorganic monovalent cations, Cs⁺ is expected to be the most feasible candidate to substitute the organic cations for two reasons.^[4] The first relates to its larger size, which satisfies the geometrical constraints of the perovskite structure to establish the continuous array of corner-sharing PbI₆ octahedra.^[5] The second arises from the superb thermal stability of CsPbX₃ (X = Cl, Br, and I) materials which may be compositionally stable even at temperatures exceeding 400 °C.^[6] So far, the efficiencies of all-inorganic Pb-halide PSCs have exceeded 20% in an n-i-p structure, which was achieved by the tetragonal (β) CsPbI₃ perovskite with an ideal band gap of 1.68 eV for photovoltaic application.^[7] Nevertheless, the imperfect Goldschmidt tolerance factor of black phase CsPbI₃ determines that their PbI₆ octahedra tends to rotate when catalyzed by H₂O, which induces a rapid phase transition to the non-photovoltaic yellow δ phase.^[8]

Several strategies were reported to address this stability issue including surface energy tuning,^[5,9] additive engineering,^[10] and interfacial passivation.^[11] However, the implementation of these strategies for improving stability is limited due to the potential impact on device performance. For instance, the amount of the additive is normally low in order to maintain crystallinity and charge mobility of the perovskite film.^[12] Similarly, the thickness of the surface passivation layer

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must be a few nanometers to maintain charge extraction from the active layers to the charge transport layers.^[13] For these reasons, highly hydrophobic hole transport materials (HTMs) have received attention, as they have the potential to block the moisture-driven phase transition of black phase CsPbI₃.^[14]

At present, state-of-the-art CsPbI₃ PSCs still use 2,2',7,7'tetrakis(N,N-di-p-methoxyphenyl-amino)-9,9'-spi-robifluorene (spiro-OMeTAD) doped with lithium bis(trifluoromethane)sulfonamide (LiTFSI), cobalt complex and 4-tert-butyl pyridine (TBP) additives as the HTM.^[7b,15] However, we observed very fast phase transitions (Figure S19) in the black phase CsPbI₃ films covered with standard doped spiro-OMeTAD layers, with these transitions being even faster than that of bare CsPbI₃ films under the same conditions. Hence, the hygroscopic nature of the dopants eliminates the advantage of the hydrophobicity of the upper hole transport layers and significantly accelerates the phase transition of CsPbI₃ leading to device degradation.^[16] Therefore, developing dopant-free, efficient, and stable HTMs is highly desirable and could lead to practical applications of inorganic perovskites in PSCs.

Herein, we described the design and synthesis of a dopantfree D- π -A type HTM, namely CI-TTIN-2F, employing a triazatruxene (TAT) as the electron-rich donor, alkylated terthiophene as π -bridges, and a fluorinated Lewis base as the electron-deficient acceptor. Benefiting from intramolecular charge transfer (ICT) and strong dipolar intermolecular interactions, CI-TTIN-2F shows excellent optoelectronic properties, ideal energy-level alignment, and good charge collection properties. In addition, joint experimental and theoretical studies suggest that CI-TTIN-2F provides multisite defect-healing effects on the surface of the CsPbI₃ films due to the presence of various heteroatoms (N, O, S, F) in the HTM structure, which effectively reduces the trap densities and charge recombination at the interface. As a dopant-free HTM for the all-inorganic CsPbI₃ PSCs, CI-TTIN-2F impressively delivers high PCEs of 15.9% on small-area cells and 11.0% on modules with a record area of 27 cm^2 . Furthermore, the device employing CI-TTIN-2F maintains over 86% of its initial performance for 1000 h after storage in ambient conditions.

Results and Discussion

Figure 1 a depicts the chemical structure of CI-TTIN-2F HTMs. The synthetic routes and the synthetic details are given in the Supplemental Information (Figures S1-S3). A planar nitrogen-containing triazatruxene (TAT) core with alkyl chains was employed as the donor (D) due to its excellent charge transporting and strong π - π stacking ability.^[17] The alkylated terthiophene conjugated arms, 3,3"dihexyl-2,2':5',2"-terthiophene, were modulated as π -bridge to increase hole-mobility properties of the molecule by increasing the double-bond character.^[18] A strong Lewis base electron-withdrawing acceptor (A) unit, 2-(5,6-difluoro-3oxo-2,3-dihydro-1*H*-inden-1-ylidene)malononitrile (IN-2F), was selected based on D-π-A type architecture to decrease the LUMO levels without causing strong steric hindrance, enhance inter/intramolecular non-covalent interactions, and increase hydrophobicity, leading to long-term stability. The expected structure of CI-TTIN-2F was confirmed by MAL-DI-TOF mass spectrometry (Figure S4). To evaluate the thermal properties of CI-TTIN-2F thermogravimetric analysis (TGA) was carried out under a nitrogen atmosphere (Figure S5). The decomposition temperature (T_{dec}) corresponding to a weight loss of 5% occurs above 300°C, indicating that CI-TTIN-2F has sufficient thermal stability for application as a HTM in PSCs. Moreover, CI-TTIN-2F is soluble in tetrachloroethane, and reasonably soluble in



Figure 1. a) Chemical structure of the CI-TTIN-2F molecule. b) UV-vis absorption spectra of CI-TTIN-2F in dilute tetrachloroethane solution. c) Spatial distribution of HOMO and LUMO orbitals in CI-TTIN-2F. C, O, N, H, F and S atoms are colored in cyan, red, blue, white, green and yellow, respectively. d) Molecular energy level alignments. e) The hole injection characteristics measured by the SCLC method of dopant-free HTMs. f) 2D GIWAXS patterns of pristine hole transport layers.

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tetrahydrofuran and dichloromethane, allowing solution processing.

The optical properties of CI-TTIN-2F were investigated by ultraviolet-visible (UV/Vis) spectroscopy in the solid state and in tetrachloroethane solution (Figure 1b). The UV-vis spectra of the CI-TTIN-2F film and solution show multimodal absorption (Figure S6). Absorption bands were observed between 350 and 500 nm, which may be attributed to the localized aromatic π - π * transition of the D- π -A structure. More importantly, an additional NIR absorption peak centered at around 660 nm indicates ultra-strong intermolecular charge transfer from the TAT electron-donating unit to the IN-2F electron-withdrawing group. Note that the position of the ICT band is slightly red-shifted as the solution transitions to the solid-state, which is characteristic of many organic semiconductors.^[19] The optical energy band gap (E_g) estimated from the onset of the absorption peak is determined to be 1.52 eV for CI-TTIN-2F, which would increase the intrinsic carrier concentration through the thermal population in the conduction band.

To investigate the highest occupied molecular orbital (HOMO) energy level (E_{HOMO}) of CI-TTIN-2F, electrochemical cyclic voltammetry (CV) was first performed with a standard three-electrode configuration (Figure S7). The material was tested in tetrahydrofuran containing 0.1 M n-Bu₄NPF₆ as a supporting electrolyte, and the oxidation potential was calibrated against an internal ferrocene standard. The E_{HOMO} value was calculated as -5.37 eV vs. vacuum for CI-TTIN-2F. The HOMO level is well aligned with the valence band energy level of the CsPbI₃ inorganic perovskite so that the photogenerated charge carriers should be efficiently transferred at the interface.^[20] The lowest unoccupied molecular orbital (LUMO) energy levels were calculated to be -3.85 eV by subtracting the optical band gap and HOMO energy level. These results are consistent with the trend estimated from density functional theory (DFT) calculations (Table S1 and S2). The molecular structures were optimized by using PBE-D3 density functional. The localization and energies of the frontier molecular levels of CI-TTIN-2F for the HOMO are delocalized primarily in the TAT core and spread along the entire molecule, while LUMO are localized in the acceptor moiety and π -bridge of the HTM (Figure 1 c). Figure 1d shows the schematic of the energy band diagram. The energy bands of the CsPbI₃ perovskite were characterized by ultraviolet photoelectron spectroscopy (UPS) and UV-vis spectroscopy as shown in Figure S8. The energy level alignment implies that the fluorinated IN-2F acceptor unit with stronger electron-withdrawing properties endows CI-TTIN-2F with much deeper HOMO energy levels by 0.40 eV compared to spiro-OMeTAD. The lower HOMO of CI-TTIN-2F is expected to ensure more efficient interfacial holetransport kinetics and improve the open-circuit voltage (V_{OC}) for CsPbI₃ PSCs.^[21]

The space-charge limited current (SCLC) measurements were carried out on the hole-only devices to evaluate the charge transport properties of CI-TTIN-2F, and the hole mobilities were extracted following the Mott-Gurney Law (Figure 1e). Encouragingly, with the D- π -A molecular structure, CI-TTIN-2F exhibits high hole mobility of $3.7 \times$ $10^{-4} \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$, significantly higher than that of spiro-OMe-TAD ($6.1 \times 10^{-5} \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$), attributed to the higher degree of conjugation and better intermolecular interactions.^[22] Grazing incidence wide-angle X-ray scattering (GIWAXS) was performed to investigate the molecular organization in the HTM films (Figure 1 f). The pole plots of the azimuth angle integrated around $q_z = 3.5 \text{ Å}^{-1}$ (Figure S9) indicate that both CI-TTIN-2F and spiro-OMeTAD possess multiple orientations in the films. Besides the face-on orientation, CI-TTIN-2F exhibits edge-on stacking that favor charge transport.^[17]

To investigate the potential interactions occurring between the CI-TTIN-2F HTM and the perovskite surface, Born Oppenheimer molecular dynamics (MD) calculations of CI-TTIN-2F deposited on top of a slab of tetragonal β-CsPbI₃ perovskite were performed (Figure 2a, movie in SI). The MD simulations show that CI-TTIN-2F can form Pb contacts of different types with the perovskite surface, specifically: i) The CN groups of the acceptor units interact with the Pb²⁺ cations of the surface, in particular forming relatively short Pb---N contacts of 2.8-3.0 Å (Figure S10), evidencing strong coordination interactions. ii) The F atoms of the acceptor can interact weakly with the surface with average Pb-F distances >4 Å (Figure S11). This is partially due to the face-on configuration of CI-TTIN-2F CN groups within 2FIN which coordinate with the perovskite surface, as malononitrile unit can rotate. Consequently, coordination of the CN group with the surface rotates the conjugated 5,6-difluoro-3-oxo-2,3dihydro-1*H*-inden-1-ylidene moiety having a quasi-planar geometry (Figures 2a), lifting the F atoms and thus elongating the Pb-F distances. However, the weak Pb-F contacts (Figure S12) contribute to more intimate interactions between the HTM and the perovskite surface. iii) The lone pairs of electrons on sulfur atoms of the π -bridges interact with the Pb²⁺ cations of the surface, which are similar to previously reported oligothiophene HTMs.^[23] The Pb-S distances around 3.7 Å are below the sum of van der Waals radii of Pb and S, evidencing coordination interactions. iv) The average Pb-O distances in the stabilized molecular conformation are about 3.0–3.5 Å, which is below or close to the sum of van-der-Waals radii of the interacting atoms. The average distances of the shortest contacts I-X (X = O, N, S, F) are longer, being about 4 Å. In the face-on configuration the hexyl units also tend to orient along the perovskite surface, which allows the sitespecific non-covalent Pb-X interactions to be maximized, since the corresponding parts of HTM can sufficiently approach the surface. Some hexyl units are unable to orient parallel to the surface due to the steric hindrance produced by other parts of CI-TTIN-2F (Figure 2a), and sometimes the hexyl chains provide steric hindrance.

To elucidate the impact of CI-TTIN-2F on the electronic structure of the interface, a representative snapshot during the MD trajectory was selected to simulate the electronic structure using PBE0 density functional (Figure 2b, c). The perovskite surface was represented by PbI₂-terminated slab, in which each Pb²⁺ ion can be considered as unsaturated, representing a defect. In the absence of CI-TTIN-2F on the perovskite surface, Pb-centered trap-like states were observed at the bottom of the conduction band of the perovskite





Figure 2. a) Molecular dynamics calculations of the CI-TTIN-2F molecule deposited on top of the slab of CsPbI₃ perovskite. b) Electronic structure of the PbI₂-terminated surface without and with CI-TTIN-2F HTMs, calculated using PBE0 density functional. c) The XPS of Pb 4f in the CsPbI₃ with and without CI-TTIN-2F films. d) The space charge limited current versus voltage of hole-only devices fabricated with spiro and CI-TTIN-2F HTMs, e) TRPL spectra and f) PL intensity maps of CsPbI₃ films on glass substrate capping with different HTMs. Scale bar, 1 μ m.

(Figure 2b). Interestingly, upon coverage of the perovskite with CI-TTIN-2F, these states disappear, with the contribution of the Pb²⁺ ions being "merged" with the continuous manifold of states at the bottom of the conduction band (Figure 2c). Moreover, the empty CI-TTIN-2F-centered states are situated below the conduction band manifold (Figure S13). Consequently, an electron from the bottom of the perovskite conduction band at the interface can migrate on the CI-TTIN-2F instead of recombining with the holes. Moreover, the participation of CI-TTIN-2F orbitals in the formation of the bottom of the conduction band favors electron delocalization over the perovskite/HTM interface, which in turn facilitates electron migration.

X-ray photoelectron spectroscopy (XPS) was carried out on the CsPbI₃ and CsPbI₃/CI-TTIN-2F films to verify the calculations. High-resolution Pb 4f spectra are shown in Figure 2d. Apart from the two main Pb $4f_{7/2}$ and Pb $4f_{5/2}$ peaks, two additional peaks around 136.9 and 141.8 eV were observed in the CsPbI₃ film attributed to the presence of the metallic Pb.^[24] The substantial metallic Pb species indicate the existence of iodide vacancies or under coordinated Pb²⁺ defects, which are likely to behave as non-radiative carrier

recombination centers and impede the solar cell performance. After depositing the CI-TTIN-2F HTM, the Pb 4f peaks shift to lower binding energies, and the two metallic Pb peaks are greatly restrained, demonstrating the strong interactions and the effective passivation of the CI-TTIN-2F HTM on the perovskite surface, in line with the MD and PDOS calculations. The influence of CI-TTIN-2F on the trap density of states of the CsPbI₃ film was further investigated by spacecharge limited current (SCLC) characterization (Figure 2e, f).^[25] The dark current density-voltage curves show three representative regions consisting of a linear ohmic, a trapfilled, and a trap-free SCLC region. The trap-filled limit voltage (V_{TFL}) is reduced to 0.88 V from 1.27 V, and the corresponding trap density of states decrease to 3.04×10^{15} from 4.38×10^{15} cm⁻³ after depositing the CI-TTIN-2F HTM. These results confirm that CI-TTIN-2F can efficiently passivate the under coordinated Pb2+ defects on the CsPbI3 perovskite surface through the interactions identified by the calculations.

Since charge extraction properties are crucial for HTMs, we examined the charge-carrier dynamics in pristine CsPbI₃ and CsPbI₃/HTM films using steady-state photoluminescence (PL) and time-resolved photoluminescence (TRPL) spectroscopy. The pristine CsPbI₃ film exhibits the strongest PL emission centered around 725 nm (Figure S14). Significant quenching was observed in the presence of the HTM, and CI-TTIN-2F has the lowest PL intensity, suggesting rapid hole extraction across the interface and a superior hole extraction ability of CI-TTIN-2F compared to spiro-OMeTAD. TRPL spectroscopy was used to delineate the carrier dynamics quantitatively (Figure 2g). The pristine CsPbI₃ film shows a longer lifetime ($\tau_1 = 2.11$ ns), which was reduced by the HTM due to the charge extraction. The CsPbI₃/CI-TTIN-2F junction presents faster hole transfer (0.69 ns) than the CsPbI₃/spiro-OMeTAD interface (0.76 ns), most likely due to the deeper valence band maximum of CI-TTIN-2F and the stronger interfacial interactions between the Pb^{2+} ion on the perovskite surface and 2FIN units of CI-TTIN-2F compared to spiro-OMeTAD.^[26] PL mapping images in Figure 2h,i reveal a similar trend where the sample with CI-TTIN-2F displays lower integrated intensity than spiro-OMeTAD. In addition, improved emission homogeneity was achieved by CI-TTIN-2F, which might result from the passivation of surface defects and the oriented face-on stacking of the HTM.

Figure 3a shows the cross-section SEM image of the complete PSCs with an n-i-p structure comprising FTO/TiO₂/ CsPbI₃/CI-TTIN-2F/Au, in which the thickness of the CI-TTIN-2F hole transport layer (HTL) was optimized at \approx 50 nm. The detailed fabrication of the devices is described in the Supporting Information, and all the HTMs were employed without any dopant if not otherwise specified. Negligible change in grain sizes and morphologies of CsPbI₃ films is observed after HTL deposition as shown in the surface-section SEM images (Figure S15). The best-performing CI-TTIN-2F-based device shows a PCE of 15.91 % under reverse scan conditions, with a short-circuit photocurrent density ($J_{\rm SC}$) of 18.82 mA cm⁻², a $V_{\rm OC}$ of 1.10 V, and a fill factor (FF) of 77.50%, and a PCE of 15.29% in forward scan condition, with a $J_{\rm SC}$ of 18.82 mA cm⁻², a $V_{\rm OC}$ of 1.06 V, and an FF of 76.90% (Figure 3b). The device performance is higher than that of the spiro-OMeTAD-based PSC with an efficiency of 11.44% under reverse scan and 8.68% in forward scan



Figure 3. a) Cross-section SEM image of the device with CI-TTIN-2F HTMs. b) The J-V curves of the CI-TTIN-2F-based champion device. c) Nyquist plots, d) Mott-Schottky plots, e) the trap density of states and f) V_{oc} as a function of light intensity of the CsPbI₃ PSCs with different HTMs. g) Photograph of the CI-TTIN-2F-based mini module. h) The J-V curve of the best-performing module. The inset is the schematic illustration of the mini-module device structure. i) PCEs of all-inorganic perovskite modules versus module area from this work and previous reports.

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(Figure S16), reflecting the lower-lying HOMO levels as well as the improved hole mobility and extraction for CI-TTIN-2F HTM. The two devices exhibit similar integrated photocurrent values (Figure S17), and the main difference in efficiency derives from the $V_{\rm OC}$ and FF. To explore the fundamental reasons, additional electrical characterization was performed on these devices. Electrochemical impedance spectroscopy (EIS) analysis reveals a smaller EIS resistance value for the CI-TTIN-2F-based device relative to the spiro-OMeTAD-based device (Figure 3 c), demonstrating its better charge-transfer behavior.^[19] The heterojunction properties at the perovskite/HTM interface were analyzed by capacitancevoltage (C-V) measurements. The $C^{-2}-V$ plots of the CI-TTIN-2F and spiro-OMeTAD-based device are depicted in Figure 3 d, following the Mott-Schottky Equation (1):

$$\frac{1}{C^2} = \frac{2}{e\varepsilon\varepsilon_0 N_{\rm D}} \left(V_{\rm bi} - V - \frac{k_{\rm B}T}{e} \right) \tag{1}$$

where C is the capacitance, $V_{\rm bi}$ refers to the built-in potential, $N_{\rm D}$ represents the charge density, and V stands for the applied voltage. $k_{\rm B}$, T, e, ε , and ε_0 are the Boltzmann constant, thermodynamic temperature, elementary charge, relative dielectric constant and vacuum permittivity, respectively. The $V_{\rm bi}$ caused by the carrier diffusion is crucial for charge injection in solar cells, which is estimated to be 1.15 V for the CI-TTIN-2F-based device, outperforming the spiro-OMe-TAD-based device (1.07 V), consistent with the decreased energy level of CI-TTIN-2F (Figure 1b). Therefore, CI-TTIN-2F extends the depleted region and enhances the driving force for carrier injection, which directly contributes to the increased $V_{\rm OC}$.^[27] Moreover, the larger slope of the Mott-Schottky plot for the CI-TTIN-2F-based device suggests a lower interfacial charge density $(3.79 \times 10^{16} \text{ cm}^{-3} \text{ vs. } 8.79 \times 10^{16} \text{ cm}^{-3} \text{ vs}^{-3} \text$ 10¹⁶ cm⁻³) and thus improved charge extraction. The higher charge accumulation at the perovskite/spiro-OMeTAD interface also explains the larger hysteresis.^[28] The trap density of states (tDOS) was also deduced from the angular frequencydependent capacitance for the fabricated devices (Figure 3e). The CI-TTIN-2F-based device demonstrates reduced trap states compared with the spiro-OMeTAD-based device over the whole defect energy region. This reduction derives from the effective surface passivation of CI-TTIN-2F, which is confirmed by the removal of trap states (Figure 2b). Accordingly, trap-assisted recombination is efficiently suppressed, as validated by the smaller ideality factor of 1.57 for the CI-TTIN-2F-based device compared with 1.92 for the spiro-OMeTAD device (Figure 3 f). Hence, CI-TTIN-2F gives rise to advantageously accelerated interface charge transfer, diminishes surface defects, and suppresses trap-assisted recombination, which reduces energy loss and thus enhances device performance, especially for $V_{\rm OC}$ and FF.

To demonstrate large-scale application of the CI-TTIN-2F HTM, inorganic CsPbI₃ perovskite modules with an active area of 27 cm² were fabricated (Figure 3 g). The modules were patterned by the P1, P2, and P3 laser-scribing method (Figure 3 h), consisting of nine sub-cells. Notably, the CsPbI₃ PSC module achieves a PCE of 10.98% employing the CI-TTIN-2F HTM, with a J_{SC} of 1.96 mA cm⁻², a V_{OC} of 9.36 V,

and an FF of 60%. The similar $V_{\rm OC}$ of each sub-cell with a single cell implies the good uniformity of the HTM layer over a large area. The reduced $J_{\rm SC}$ and FF can be attributed to the greater series resistance because of the longer carrier transport paths and the device nonuniformities. Moreover, the perovskite solar module with the CI-TTIN-2F HTM shows a steady output PCE of 10.80% under the AM 1.5 G illumination for 250 s, which is consistent with that calculated from J-V scanning (Figure S18). The performance of the reported inorganic perovskite modules versus device area are summarized in Figure 3i.^[29] The efficiency of the CI-TTIN-2F-based module is comparable with that of previous reports while showing the largest module area.

Moisture-induced phase transitions of the perovskite film are a critical degradation path for all-inorganic PSCs. To examine the effects of HTL incorporation on the phase stability of perovskite films, we tracked the absorbance evolution of CsPbI₃ films under a controlled relative humidity (RH) of $\approx 50\%$, and photographs of the perovskite films stored for different times are shown in Figure 4a and Figure S19. Remarkably, the CsPbI₃ film covered with doped spiro-OMeTAD is bleached within only 20 min, which is even more rapid than that of the bare CsPbI₃ film (100 min). The degradation onset is significantly retarded to 180 min when the hydrophilic dopants were removed from the spiro-OMeTAD HTL composition. Specially, no obvious change in color is observed from the CsPbI₃ film with the CI-TTIN-2F HTM after aging for 480 min. In addition, the long-term phase stability of the films under ambient conditions with an RH of $\approx 20\%$ was further investigated by X-ray diffraction (XRD). All the fresh CsPbI₃ films show a typical black phase with two main specific peaks located at 14.6° and 29.2° which are assigned to (110) and (220) planes, respectively (Figure 4b). The CsPbI₃ films with various HTLs behave differently after aging for 360 h (Figure 4c). A characteristic peak assigned to the δ phase at 10.2° was observed in the bare CsPbI₃ film after aging.^[30] Significantly, the upper spiro-OMeTAD HTL with and without dopants accelerate and decelerate, respectively, the phase transition of the CsPbI₃ perovskite structure to the yellow δ phase. Thus, these hygroscopic dopants not only eliminate the advantage of hydrophobicity of the upper spiro-OMeTAD layer but also further accelerate the phase transition of the CsPbI₃ films. When coated with the dopant-free CI-TTIN-2F HTL, the CsPbI₃ film remains in the black phase and is almost identical to the fresh film, indicating that CI-TTIN-2F is capable of protecting the metastable perovskite films from moisture penetration. To probe surface property of the HTLs, contact angles of water droplets on the CsPbI₃/HTLs samples were measured (Figure 4d). The CI-TTIN-2F HTL showed a much higher contact angle of $\approx 99^{\circ}$ than those of spiro-OMeTAD with ($\approx 75^\circ$) and without dopants ($\approx 88^\circ$), and the improved hydrophobicity can be in part attributed to the introduced fluorinated 1,1-dicyanomethylene-3-indanone acceptor units in the D- π -A structure.

The stability of unencapsulated CsPbI₃ PSCs employing different HTMs was monitored in an ambient atmosphere with an RH of ≈ 20 % for 1000 hours. Figure 4e presents the normalized PCEs of the devices as a function of storage time,



Figure 4. a) Photos of the CsPbl₃ films with and without different HTMs before and after exposure to RH of \approx 50% for 480 min. XRD pattern of the CsPbl₃ films with and without different HTMs (b) before and (c) after exposure to RH of \approx 20% for 360 h. d) Water contact angles of different HTLs. e) Stability under ambient (\approx 20% RH) for the unencapsulated CsPbl₃ PSCs with different HTMs. f) Stability of encapsulated CI-TTIN-2F-based device in air (\approx 50% RH), the inset is the illustration of the cross sections of the respective encapsulation schemes. g) Operational stability of the CsPbl₃ device with CI-TTIN-2F HTMs under a constant one sun illumination at the maximum power point without a UV filter in nitrogen atmosphere.

and the device performances were periodically measured in ambient air. For comparison, we also fabricated CsPbI3 devices using spiro-OMeTAD with conventional dopants and the best-performing device exhibited a PCE of 17.77 % with a $J_{\rm SC}$ of 19.53 mA cm⁻², a $V_{\rm OC}$ of 1.12 V, and an FF of 78.70% (Figure S20). However, the ambient stability monitored under ambient conditions reveals that the doped spiro-OMeTAD-based device suffers a sharp drop up to < 10% of the initial PCE within only 100 h due to the moisture-sensitive dopants. Furthermore, the dopant-free spiro-OMeTADbased device retains 66% of the pristine performance after 1000 hours. Encouragingly, the CI-TTIN-2F-based device maintains 86% of its initial PCE after 1000 h exposure. To further enhance device stability, we encapsulated the CI-TTIN-2F-based device and a cross-section of the encapsulation Scheme is illustrated in the inset in Figure 4 f. Successfully, the device displays improved stability, and the efficiency drops by only 5% after storing under a controlled \approx 50% RH for over 1500 h. Moreover, we examined the operational stability of the unencapsulated CI-TTIN-2F-based device under a constant one sun illumination (AM1.5G) at 25 °C in a nitrogen atmosphere at the maximum powerpoint (Figure 4g). The CsPbI₃ PSC with CI-TTIN-2F maintains over 80% of the initial PCE under light soaking for 800 h, demonstrating a superior photochemical stability based on this new CI-TTIN-2F HTM.

Conclusion

We synthesized a new dopant-free HTM named CI-TTIN-2F with a D- π -A molecular configuration. The suitable optoelectronic properties and energy-level alignment endow CI-TTIN-2F with excellent charge collection properties. In addition, joint experimental and theoretical studies suggest that CI-TTIN-2F has the capacity for multisite passivation effects on defective CsPbI₃ surfaces due to the presence of various heteroatoms. As a result, all-inorganic CsPbI₃ PSCs with dopant-free CI-TTIN-2F HTM demonstrate a high PCE of 15.9%, along with 86% efficiency retention after

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1000 hours under ambient conditions. These results indicate an excellent compatibility of CI-TTIN-2F with all-inorganic PSCs. Notably, the largest all-inorganic perovskite solar module was fabricated using the CI-TTIN-2F HTM, and exhibited an efficiency of 11.0 % with an area of 27 cm². This study provides a new design strategy toward efficient dopantfree HTMs with multisite passivation effects that stabilize allinorganic PSCs to facilitate future scale-up.

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Conflict of Interest

The authors declare no conflict of interest.

Keywords: defect passivation · dopant-free · hole transport materials · inorganic perovskites · perovskite solar cells

- a) J. J. Yoo, G. Seo, M. R. Chua, T. G. Park, Y. Lu, F. Rotermund, Y.-K. Kim, C. S. Moon, N. J. Jeon, J.-P. Correa-Baena, *Nature* 2021, 590, 587–593; b) J. Jeong, M. Kim, J. Seo, H. Lu, P. Ahlawat, A. Mishra, Y. Yang, M. A. Hope, F. T. Eickemeyer, M. Kim, *Nature* 2021, 592, 381–385; c) NREL, Best Research-Cell Efficiences, https://www.nrel.gov/pv/cell-efficiency.html (accessed: May 2021).
- [2] a) J. Zhang, G. Hodes, Z. Jin, S. Liu, Angew. Chem. Int. Ed. 2019, 58, 15596–15618; Angew. Chem. 2019, 131, 15742–15765; b) Z. Yao, W. Zhao, S. F. Liu, J. Mater. Chem. A 2021, 9, 11124–11144.
- [3] J. Tian, Q. Xue, Q. Yao, N. Li, C. J. Brabec, H. L. Yip, Adv. Energy Mater. 2020, 10, 2000183.
- [4] X. Sun, Z. Shao, Y. Rao, H. Meng, C. Gao, C. Chen, D. Liu, P. Lv, Z. Li, X. Wang, Adv. Energy Mater. 2021, 11, 2002754.
- [5] C. Liu, Y. Yang, O. A. Syzgantseva, Y. Ding, M. A. Syzgantseva, X. Zhang, A. M. Asiri, S. Dai, M. K. Nazeeruddin, *Adv. Mater.* 2020, 32, 2002632.
- [6] R. J. Sutton, M. R. Filip, A. A. Haghighirad, N. Sakai, B. Wenger, F. Giustino, H. J. Snaith, ACS Energy Lett. 2018, 3, 1787–1794.

- [7] a) Y. Wang, M. I. Dar, L. K. Ono, T. Zhang, M. Kan, Y. Li, L. Zhang, X. Wang, Y. Yang, X. Gao, *Science* 2019, *365*, 591–595;
 b) C. C. Boyd, R. Cheacharoen, T. Leijtens, M. D. McGehee, *Chem. Rev.* 2019, *119*, 3418–3451.
- [8] Z. Li, F. Zhou, Q. Wang, L. Ding, Z. Jin, *Nano Energy* 2020, 71, 104634.
- [9] a) Q. Wang, X. Zheng, Y. Deng, J. Zhao, Z. Chen, J. Huang, *Joule* 2017, 1, 371–382; b) J. Xi, C. Piao, J. Byeon, J. Yoon, Z. Wu, M. Choi, *Adv. Energy Mater.* 2019, 9, 1901787; c) T. Zhang, M. I. Dar, G. Li, F. Xu, N. Guo, M. Grätzel, Y. Zhao, *Sci. Adv.* 2017, 3, e1700841.
- [10] T. Wu, Y. Wang, Z. Dai, D. Cui, T. Wang, X. Meng, E. Bi, X. Yang, L. Han, Adv. Mater. 2019, 31, 1900605.
- [11] a) K. Wang, Q. Tian, G. Zhao, J. Wen, J. Huang, C. Gao, Z. Xu, Y. Liu, L. Liang, L. Meng, *Cell Rep. Phys. Sci.* **2020**, *1*, 100180; b) T. Liu, J. Zhang, M. Qin, X. Wu, F. Li, X. Lu, Z. Zhu, A. K. Y. Jen, *Adv. Funct. Mater.* **2021**, *31*, 2009515; c) J. Wang, J. Zhang, Y. Zhou, H. Liu, Q. Xue, X. Li, C.-C. Chueh, H.-L. Yip, Z. Zhu, A. K. Jen, *Nat. Commun.* **2020**, *11*, 1–9.
- [12] C. Liu, Y. Yang, X. Liu, Y. Ding, Z. Arain, X. Li, Y. Li, Z. Zhou, S. Dai, M. K. Nazeeruddin, *J. Mater. Chem. A* **2020**, *8*, 10226– 10232.
- [13] C. Yan, Z. Li, Y. Sun, J. Zhao, X. Huang, J. Yang, Z. Ci, L. Ding, Z. Jin, J. Mater. Chem. A 2020, 8, 10346–10353.
- [14] a) T. Miyasaka, A. Kulkarni, G. M. Kim, S. Öz, A. K. Jena, Adv. Energy Mater. 2020, 10, 1902500; b) Z. Guo, A. K. Jena, I. Takei, G. M. Kim, M. A. Kamarudin, Y. Sanehira, A. Ishii, Y. Numata, S. Hayase, T. Miyasaka, J. Am. Chem. Soc. 2020, 142, 9725 – 9734; c) M. Cheng, C. Zuo, Y. Wu, Z. Li, B. Xu, Y. Hua, L. Ding, Sci. Bull. 2020, 65, 1237; d) X. Ding, H. Wang, C. Chen, H. Li, Y. Tian, Q. Li, C. Wu, L. Ding, X. Yang, M. Cheng, Chem. Eng. J. 2021, 410, 128328.
- [15] X. Chang, J. Fang, Y. Fan, T. Luo, H. Su, Y. Zhang, J. Lu, L. Tsetseris, T. D. Anthopoulos, S. Liu, *Adv. Mater.* **2020**, *32*, 2001243.
- [16] a) W. Xiang, W. Tress, Adv. Mater. 2019, 31, 1902851; b) S. S. Mali, J. V. Patil, J. A. Steele, S. R. Rondiya, N. Y. Dzade, C. K. Hong, ACS Energy Lett. 2021, 6, 778–788; c) S. Öz, A. K. Jena, A. Kulkarni, K. Mouri, T. Yokoyama, I. Takei, F. Ünlü, S. Mathur, T. Miyasaka, ACS Energy Lett. 2020, 5, 1292–1299.
- [17] C. Igci, S. Paek, K. Rakstys, H. Kanda, N. Shibayama, V. Jankauskas, C. Roldán-Carmona, H. Kim, A. M. Asiri, M. K. Nazeeruddin, *Solar RRL* 2020, *4*, 2000173.
- [18] K. Rakstys, S. Paek, P. Gao, P. Gratia, T. Marszalek, G. Grancini, K. T. Cho, K. Genevicius, V. Jankauskas, W. Pisula, *J. Mater. Chem. A* 2017, *5*, 7811–7815.
- [19] A. Kojima, K. Teshima, Y. Shirai, T. Miyasaka, J. Am. Chem. Soc. 2009, 131, 6050–6051.
- [20] a) C. Liu, L. Zhang, Y. Li, X. Zhou, S. She, X. Wang, Y. Tian, A. K. Jen, B. Xu, *Adv. Funct. Mater.* 2020, *30*, 1908462; b) Y. Ou, A. Sun, H. Li, T. Wu, D. Zhang, P. Xu, R. Zhao, L. Zhu, R. Wang, B. Xu, Y. Hua, L. Ding, *Mater. Chem. Front.* 2021, *5*, 876; c) D. Zhang, T. Wu, P. Xu, Y. Ou, A. Sun, H. Ma, B. Cui, H. Sun, L. Ding, Y. Hua, *J. Mater. Chem. A* 2019, *7*, 10319.
- [21] T. Niu, W. Zhu, Y. Zhang, Q. Xue, X. Jiao, Z. Wang, Y.-M. Xie, P. Li, R. Chen, F. Huang, *Joule* **2021**, *5*, 249–269.
- [22] S. Paek, P. Qin, Y. Lee, K. T. Cho, P. Gao, G. Grancini, E. Oveisi, P. Gratia, K. Rakstys, S. A. Al-Muhtaseb, *Adv. Mater.* 2017, 29, 1606555.
- [23] Q. Xiao, T. Sakurai, T. Fukino, K. Akaike, Y. Honsho, A. Saeki, S. Seki, K. Kato, M. Takata, T. Aida, *J. Am. Chem. Soc.* 2013, 135, 18268–18271.
- [24] D. Bi, C. Yi, J. Luo, J.-D. Décoppet, F. Zhang, S. M. Zakeeruddin, X. Li, A. Hagfeldt, M. Grätzel, *Nat. Energy* **2016**, *1*, 16142.
- [25] a) Y. Yang, C. Liu, A. Mahata, M. Li, C. Roldán-Carmona, Y. Ding, Z. Arain, W. Xu, Y. Yang, P. A. Schouwink, A. Züttel, F. De Angelis, S. Dai, M. K. Nazeeruddin, *Energy Environ. Sci.*

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2020, *13*, 3093–3101; b) Y. Yang, C. Liu, O. A. Syzgantseva, M. A. Syzgantseva, S. Ma, Y. Ding, M. Cai, X. Liu, S. Dai, M. K. Nazeeruddin, *Adv. Energy Mater.* **2021**, *11*, 2002966.

- [26] N. Arora, M. I. Dar, A. Hinderhofer, N. Pellet, F. Schreiber, S. M. Zakeeruddin, M. Grätzel, *Science* 2017, 358, 768–771.
- [27] M. Du, X. Zhu, L. Wang, H. Wang, J. Feng, X. Jiang, Y. Cao, Y. Sun, L. Duan, Y. Jiao, Adv. Mater. 2020, 32, 2004979.
- [28] Y. Liu, M. Bag, L. A. Renna, Z. A. Page, P. Kim, T. Emrick, D. Venkataraman, T. P. Russell, *Adv. Energy Mater.* **2016**, *6*, 1501606.
- [29] a) C. Liu, Y. Yang, C. Zhang, S. Wu, L. Wei, F. Guo, G. M. Arumugam, J. Hu, X. Liu, J. Lin, Adv. Mater. 2020, 32, 1907361;

b) R. Chen, Y. Hui, B. Wu, Y. Wang, X. Huang, Z. Xu, P. Ruan, W. Zhang, F. Cheng, W. Zhang, J. Mater. Chem. A 2020, 8, 9597– 9606.

[30] C. Liu, Y. Yang, X. Xia, Y. Ding, Z. Arain, S. An, X. Liu, R. C. Cristina, S. Dai, M. K. Nazeeruddin, *Adv. Energy Mater.* 2020, *10*, 1903751.

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